



Highly enhanced electron injection in organic light-emitting diodes with an n-type semiconducting MnO₂ layer

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ABSTRACT

Highly enhanced electron injection is demonstrated with a thin manganese dioxide (MnO₂) electron injection layer (EIL) in Alq₃-based organic light-emitting diodes. Insertion of the MnO₂ EIL between the Al cathode and Alq₃ results in highly improved device characteristics. *In situ* photoelectron spectroscopy shows remarkable reduction of the electron injection barrier without significant chemical reactions between Alq₃ and MnO₂, which could induce Alq₃ destruction. The reduction of the electron injection barrier is due to the n-type doping effect, and the lack of strong interfacial reaction is advantageous with regards to more efficient electron injection than a conventional LiF EIL. These properties render the MnO₂ a potential EIL.

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1. Introduction

Organic light-emitting diodes (OLEDs) have been remarkably attractive for the last two decades due to their unique advantages, such as simple fabrication processes, wide viewing angles, light weight, and mechanical flexibility [1]. However, high charge-injection barrier from the electrode to organic semiconducting materials is one of the inveterate obstacles to design highly efficient OLEDs. Lowering the electron injection barrier in OLEDs especially is a prerequisite to balancing electrons with holes within an emission layer to maximize light emission because most organic semiconducting materials have lower electron mobility than that of hole. It has been reported that the insertion of an appropriate electron injection layer (EIL) reduces the electron injection barrier efficiently.

Alkali metals and alkali metal halides have been used as a conventional EIL through the n-type doping effect on an electron transport layer (ETL), such as Alq₃ [2,3]. However, alkali metals are not tractable as an evaporation source and are destructively reactive to the organic molecules. On the other hand, alkali metal halides often show variable performance depending upon the cathode choice [4]. Recently, several alternative EILs, such as alkali metal carbonates, nitride and quinolates [5–11] have been studied and show good electron injection performances. Metal oxides such as ZnO, TiO₂ and ZrO₂ are also candidates for alternative EILs [12–14]. Recently, Luo et al. reported highly efficient OLEDs with insulating manganese monoxide (MnO) EIL [15]. However, its detailed working mechanism is not yet understood. Furthermore, the semiconducting phase of metal oxide would be more suitable for device performance than the insulating phase.

In this paper, we propose that semiconducting manganese dioxide (MnO₂) could be a potential candidate as a

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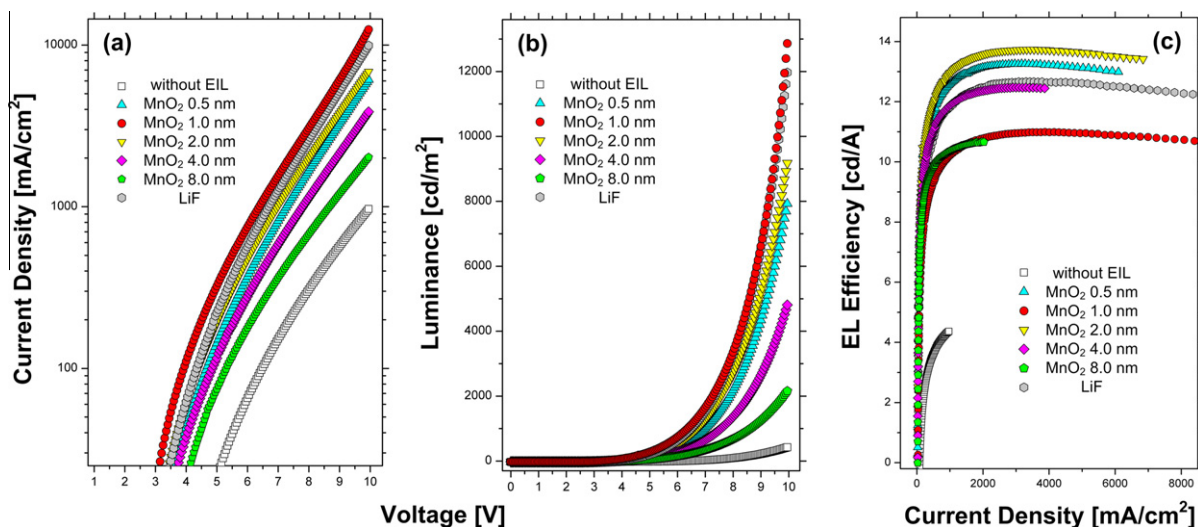


Fig. 1. Measured (a) current density–voltage, (b) luminance–voltage characteristics and (c) EL efficiency of OLEDs with different thicknesses (0, 0.5, 1.0, 2.0, 4.0 and 8.0 nm) of an MnO₂ electron injection layer. The device characteristics with a conventional LiF EIL are also depicted to show the superiority of a MnO₂ EIL.

highly efficient EIL. First, we observed conspicuous enhancements of electron injections with a MnO₂ EIL through OLEDs fabrication including current density–voltage–luminance (J – V – L) measurements. We also discovered the optimum thickness of a MnO₂ layer and compared that with a conventional LiF EIL. To elucidate the origin of the electron injection enhancements, we carried out *in situ* photoelectron spectroscopy (PES) experiments. This allowed us to understand the interfacial electronic structures and chemistry which play a crucial role in device performance.

2. Experimental

We fabricated OLEDs with the structure of Al (100 nm)/MnO₂ (0, 0.5, 1.0, 2.0, 4.0, 8.0 nm)/Alq₃ (50 nm)/NPB (50 nm)/ITO. Organic materials and MnO₂ (Sigma Aldrich, 99.99+%) were successively deposited onto the ITO patterned glass substrate by thermal evaporation in an ultra-high vacuum (UHV) chamber below 5×10^{-8} and 2×10^{-7} Torr at the rates of 0.1 nm/s and 0.01 nm/s, respectively. Devices were then completed by depositing the Al cathode on the sample with the rate of >0.1 nm/s. All the deposition rates and thicknesses were monitored by a quartz crystal microbalance. Active areas of the devices are 0.04 cm². J – V – L characteristics were measured under a dry nitrogen atmosphere using a Keithely 237 and 2400 source measure unit with a photodiode calibrated by a PR650 spectrophotometer. To show the superior electron injection of MnO₂ EIL, OLED with LiF EIL is also fabricated with its optimized thickness (0.1 nm) in our preparation condition.

In situ PES experiments were performed using a PSP RESOLVE 120 spectrometer in an analysis chamber which is directly connected with a preparation chamber. An ultraviolet (He I, 21.22 eV) and a non-monochromatized X-ray (Mg K α , 1253.6 eV) radiations were used as excitation light

sources. To obtain the secondary electron cutoff (SEC), a sample bias of -10 V was applied in normal emission geometry. We used the same interface formation to that of device fabrication because the electronic structures are significantly varied by their deposition sequence between organic and metal electrode [16]. To investigate the effects of the MnO₂ layer insertion, we prepared Al/Alq₃ and Al/MnO₂/Alq₃ samples and compared their electronic structures. In both cases, ITO coated glasses were used as substrates after checking if they all had the same work function. The analysis and preparation chamber were maintained at pressures of 3×10^{-9} and 2×10^{-7} Torr, respectively. The deposition rates were cross checked by a calibrated thickness monitor and by the attenuation of core level intensities in X-ray photoelectron spectroscopy. The interfacial electronic structures of Al/LiF/Alq₃ were referred from well-known reports [4,17].

3. Results and discussion

Fig. 1a and b depicts measured J – V – L characteristics. Device performances were highly improved by the insertion of the MnO₂ layer. The optimized thickness of a MnO₂ layer was observed at 1.0 nm and device performance decreased gradually as the MnO₂ layer thickened due to incremental series resistance. However, the device with a relatively thick MnO₂ layer (8.0 nm) still showed better performance than the reference device without MnO₂ insertion. This thickness-independent stability in the device performance originates from the semiconducting nature of MnO₂, which is clear contrast to common EILs such as insulating alkali metal halides [18,19]. As the MnO₂ layer of 1.0 nm is inserted, both current density and luminance increase by more than one order of magnitude at 10 V compared to those of the reference device. Turn on voltage, which is defined as the voltage giving 1 cd/m² luminance, was also significantly reduced from 6.4 V to

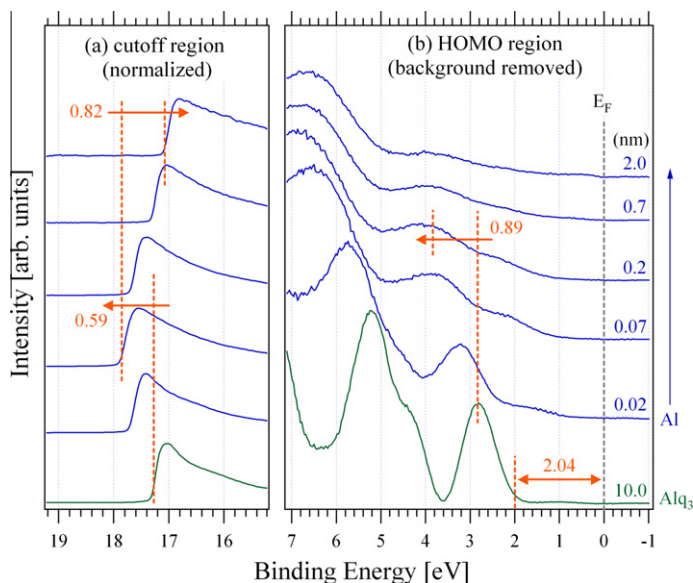


Fig. 2. Measured UPS spectra of (a) the secondary electron cutoff and (b) the HOMO region of Al (0.02, 0.07, 0.2, 0.7 and 2.0 nm)/Alq₃ (10.0 nm). Each spectrum is (a) normalized and (b) background removed for easy comparison, respectively.

3.3 V as the 1.0 nm-MnO₂ was inserted. Furthermore, OLED with 1.0 nm-MnO₂ shows more than 25% higher current density than that of LiF EIL at 10 V. However, as shown in the Fig. 1c, the device with 1.0 nm-MnO₂ interlayer has poorer electroluminescence (EL) efficiency than those with other thicknesses or LiF. Since it clearly shows significant improvement in electron injection from the Al cathode, this low efficiency must be due to the imbalance between electrons and holes within the light emission layer by the excessive electron injection. If one adopted an appropriate hole injection layer too, such as MoO₃ or WO₃ [20,21], EL efficiency could be optimized with an MnO₂ layer of 1.0 nm thickness.

Fig. 2 shows the measured UPS spectra of (a) SEC and (b) the highest occupied molecular orbital (HOMO) region of Al (0.02, 0.07, 0.2, 0.7 and 2.0 nm)/Alq₃ (10.0 nm). The onset of HOMO level of 10.0 nm-Alq₃ is seen at 2.04 eV below the Fermi level. As Al is deposited on Alq₃, SEC shifts to the higher binding energy side until 0.07 nm-Al is deposited. Simultaneously, the HOMO level shifts gradually to the higher binding energy side and the gap state emerges at 2 eV below the Fermi level. As a result, Alq₃ is weakly doped in n-type due to the chemical interaction between Al and Alq₃ without EIL deposition on Alq₃. This interaction has already reported and our results accord well [22]. The total HOMO level shift of Alq₃ is 0.89 eV and the total SEC shift (Δ SEC) is 0.59 eV. After 0.07 nm-Al deposition, SEC shifts to the lower binding energy side, approaching to the work function of Al. Meanwhile, the HOMO level of Alq₃ does not shift any more after 0.2 nm-Al deposition and its intensity is gradually attenuated. Finally, the Fermi level of Al is seen with 2.0 nm-Al deposition.

Fig. 3 displays the measured UPS spectra of (a) SEC and (b) HOMO region of Al (0.02, 0.07, 0.3, 1.0 and 4.0 nm)/MnO₂ (0.5 and 1.0 nm)/Alq₃ (10.0 nm). The onset of HOMO level of 10.0 nm-Alq₃ is seen at 2.06 eV from the Fermi level.

As MnO₂ is deposited on Alq₃, the SEC shift to the higher binding energy side is notably larger than that of Al/Alq₃. After 1.0 nm deposition of MnO₂, the SEC shifts by 1.31 eV to the higher binding energy side. Simultaneously, the HOMO level of Alq₃ also shifts by 1.43 eV to the higher binding energy side, which is much larger than that without MnO₂. This implies that the MnO₂ EIL could work efficiently irrespective of the cathode choice unlike LiF/Al vs. LiF/Ag [23,24], where the large HOMO level shift occurs without Al deposition (This will be proven in following device results.) In addition, the amount of HOMO level shift upon the MnO₂ deposition is larger than the reports for Al/LiF/Alq₃, thus it accords well with our device results [4,17,25]. This shift is comparable with that of n-type doping using low work function metals [26]. Furthermore, a broad and weak gap state appears at the high energy position by 1.4 eV apart from the HOMO [inset in (b)]. However, since this gap state becomes visible only after background subtraction and fitting, the interaction between Alq₃ and MnO₂ would be very weak, compared with common interaction features between alkali metals and Alq₃ [27]. Comparing the MnO₂ 1.0 nm spectra with that of 0.5 nm, the intensity of the gap state does not noticeably increase. In addition, the valence electronic structure of the pristine Alq₃ is well conserved except for the appearance of a weak gap state upon the MnO₂ deposition. This implies that the n-type doping, electron transfer from MnO₂ to Alq₃, occurs effectively without significant molecular destruction. Ding et al. reported that n-type doping with Cs atoms on organic materials occurred through the following two sequential interactions: (1) the LUMO level shift to the Fermi level at low coverage and (2) the gap state appearance from the LUMO level filling without further energy position shift at high coverage [28,29]. While the interaction (1) reduces the electron injection barrier, the interaction (2) would not contribute to efficient

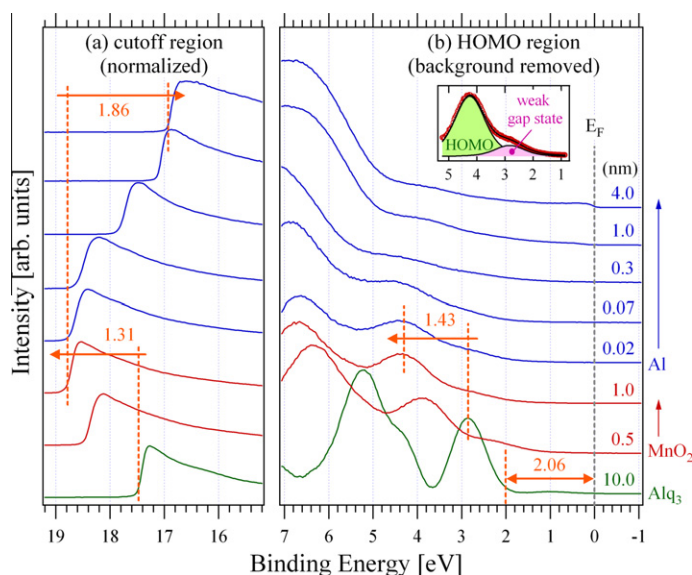


Fig. 3. Measured UPS spectra of (a) the secondary electron cutoff and (b) the HOMO region of Al (0.02, 0.07, 0.3, 1.0 and 4.0 nm)/MnO₂ (0.5 and 1.0 nm)/Alq₃ (10.0 nm). Each spectrum is (a) normalized and (b) background removed for easy comparison, respectively. Magnified HOMO region at MnO₂ 1.0 nm is shown in the inset to show the weak broad gap state with peak fitting.

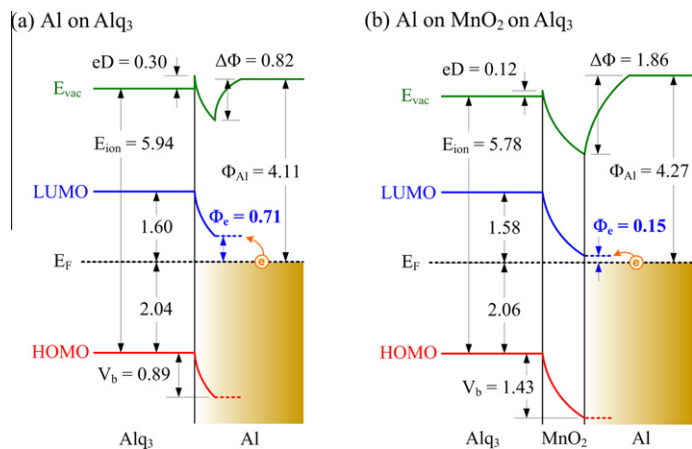


Fig. 4. Energy level diagrams of (a) Al/Alq₃ and (b) Al/MnO₂/Alq₃. Φ_e is the electron injection barrier and E_{ion} is the ionization energy of Alq₃. The band bending (V_b) was evaluated from the shift of the HOMO peak during deposition. Gap states are omitted for brevity.

electron injection anymore and even deteriorates the device performance, because the density of states of the pristine Alq₃ LUMO decreases due to the gap state filling. In general, MnO₂ is known as an n-type semiconductor [30]. Furthermore, thermally evaporated transition metal oxides contain substantial amount of oxygen vacancies, which gives them stronger n-type characteristics [31]. Therefore, Alq₃ could receive electrons from thermally deposited MnO₂ which have enough n-type carriers, i.e. electrons. However, it is expected that the electron transfer from MnO₂ is weaker than that of alkali metals, thus the interaction (1) would be the dominant effect with the MnO₂ EIL. This assumption accords well with the behavior observed in PES spectra where there were large HOMO level shifts indicating significant interaction (1) and a weak gap state

indicating insignificant interaction (2). Upon deposition of the Al layer (blue spectra) on the MnO₂ deposited Alq₃, the SEC moves back to lower binding energy side while the HOMO level hardly shifts. The Fermi level of Al is first seen at the Al 1.0 nm deposition step and clearly seen at 4.0 nm. The work function approaches 4.27 eV as the Al layer thickens, which corresponds well with the work function of the Al metal.

Combining all information from the measured spectra, we drew the energy level diagrams of Al/Alq₃ (spectra not shown) and Al/MnO₂/Alq₃ in Fig. 4. The LUMO level of Alq₃ was estimated from the transport band gap of 3.64 eV obtained from the previous report using PES and inverse photoelectron spectroscopy [32]. The ionization energy (E_{ion}) of Alq₃ and the bulk Al work function are both

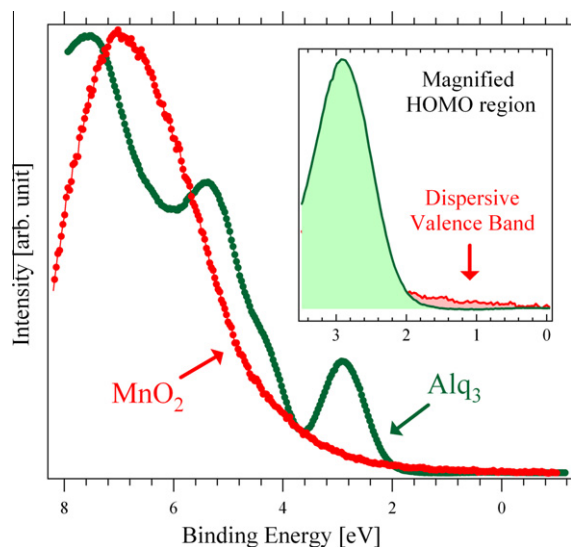


Fig. 5. Measured valence/HOMO region of neat MnO₂ film (5 nm, red mark and line) and Alq₃ film (10 nm, green mark and line) on ITO. To compare them easily, magnified HOMO region is also depicted in the inset. The dispersive valence band of MnO₂ is clearly seen above the HOMO of Alq₃. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

within acceptable ranges with regard to the resolution of our system. Band bending (V_b) is determined by the value of the HOMO peak shift and the interfacial dipole (eD) is evaluated by subtracting V_b from ΔSEC . In Fig. 4a, the SEC and HOMO level shift to higher binding energy side upon Al deposition because of the interaction between Al and Alq₃, as shown in PES spectra. Finally, the SEC reaches the work function of Al metal (4.11 eV). As a result, the electron injection barrier (Φ_e) from Al to Alq₃ is measured as 0.71 eV. However, in Fig. 4b, the MnO₂ layer induces a much larger HOMO level shift than the direct Al/Alq₃ inter-

face. No further HOMO shift is observed by subsequent Al deposition, which suggests that the MnO₂ could prevent Al–Alq₃ reactions and would contribute to the device stability. Consequently, Φ_e is dramatically reduced to 0.15 eV, and thus electrons are easily injected from the Al cathode to Alq₃. This is the origin of the electron injection and device performance improvements with a MnO₂ EIL.

To elucidate the origin of the electron transfer from MnO₂ to Alq₃, thus the n-type doping, we measured the valence spectrum of neat MnO₂ thin film and compare it with the HOMO region of Alq₃. We drew the spectra in Fig. 5 with the common vacuum level of neat MnO₂ and Alq₃ film to see the relative level position before the contact formation between MnO₂ and Alq₃. Both spectra were normalized, and backgrounds and He I _{β} satellites were removed to avoid any confusion. Fig. 5 shows the valence region spectra of MnO₂ (5 nm, red marks and line) and Alq₃ (10 nm, green marks and line) on ITO. While the strong feature of MnO₂ is seen in the range of 5–8 eV, the end tail (emission features at <5 eV) of the dispersive valence band is widely spread over several eVs. This dispersive band is well coincided with the previous XPS and theoretical study [33]. To compare the dispersive valence band of MnO₂ with the HOMO of Alq₃ more clearly, the magnified HOMO region of both spectra are shown in the inset of Fig. 5. The definite dispersive valence band of MnO₂ exists, although its intensity is relatively weak (Generally, in the case of metal oxide film, the end tail of the dispersive energy band is very weak and even it is not observable in PES spectra [34,35].) This density of occupied state could originate from the oxygen vacancies by thermal evaporation, as we aforementioned [31]. The density of occupied state of MnO₂ is broadly distributed far above the HOMO onset of Alq₃ (~ 2 eV), thus the electrons in this state could easily transfer to the LUMO of Alq₃. This is the origin of the n-type doping properties of MnO₂.

To prove the independence of the cathode choice, we measured the J – V characteristics of electron-only devices

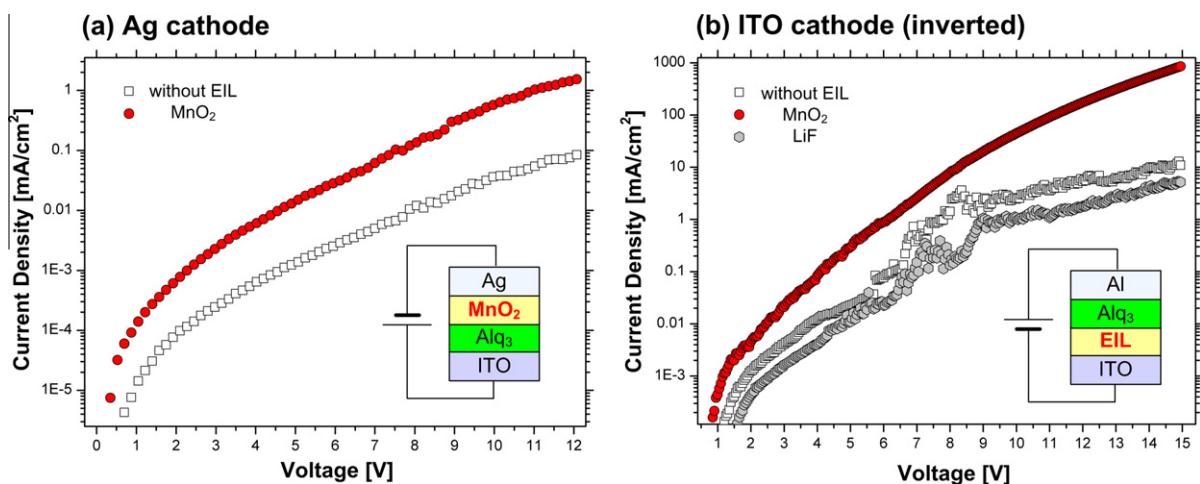


Fig. 6. Measured J – V characteristics of electron-only devices with (a) Ag cathode (Ag/MnO₂/Alq₃/ITO) and (b) the inverted structure (Al/Alq₃/EIL/ITO). Schematic configuration of the device is shown in the inset. The MnO₂ devices show much better performance with any cathode–MnO₂ combination than that without EIL.

with two different cathodes in combination with MnO₂ EIL. First, we fabricated the electron-only devices with Ag cathode, which has the device structure of Ag (35 nm)/MnO₂ (1 nm or without)/Alq₃ (300 nm)/ITO. Fig. 6a depicts the measured *J*–*V* characteristics. At the applied bias of 12 V, MnO₂ device shows higher current density more than one order of magnitude compared to the current density of the device without EIL. This is the identical result, as the *J*–*V* characteristics of Al/MnO₂ OLEDs, thus MnO₂ EIL works efficiently with Ag cathode, unlike LiF [4,23,24]. In addition, we measured the *J*–*V* characteristics of the inverted device structures, which consist of Al (100 nm)/Alq₃ (75 nm)/[MnO₂ (1.0 nm)] or [LiF (0.1 nm)] or [without EIL]/ITO. Other conditions are identical as those of our normal OLED structure. In this inverted configuration, ITO was used as a cathode. Measured *J*–*V* results are shown in Fig. 6b. At 15 V, MnO₂ device shows more than 80 times higher current density than that without EIL. However, LiF EIL does not enhance the *J*–*V* characteristics at all, even deteriorates, implying that LiF just acts as a resistance. Therefore, MnO₂ is more versatile to be used as an EIL than LiF. In addition, however, we remark the extrinsic factors, such as growth mode of MnO₂ and LiF on ITO, could also highly affect on the *J*–*V* characteristics.

4. Conclusion

In summary, we investigated significant enhancements of electron injection by inserting the MnO₂ EIL in OLEDs. Device characteristics increase remarkably with MnO₂ EIL and the interfacial electronic structures of Al/Alq₃ and Al/MnO₂/Alq₃ show the prominent reduction of the electron injection barrier from Al to Alq₃, which is comparable or even smaller than that of the conventional EIL, such as LiF or reactive metals. In addition, MnO₂ does not show strong interaction with Alq₃ unlike alkali metals and their halides, which induce molecular distortion/destruction, and instead merely dopes the Alq₃ layer efficiently. The reduction of the electron injection barrier originates from the electron doping effect of MnO₂ having a highly n-type semiconducting nature. Furthermore, efficient electron injection is observed even when MnO₂ is adopted in the device with Ag cathode or inverted structure. These advantages render MnO₂ as a potential alternative for highly efficient EIL.

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References

- [1] C.W. Tang, S.A. VanSlyke, *Appl. Phys. Lett.* 51 (1987) 913.
- [2] T. Oyamada, C. Maeda, H. Sasabe, C. Adachi, *Jap. J. Appl. Phys.* 42 (2003) 1535.
- [3] L.S. Hung, C.W. Tang, M.G. Mason, *Appl. Phys. Lett.* 70 (1997) 152.
- [4] C.-I. Wu, G.-R. Lee, T.-W. Pi, *Appl. Phys. Lett.* 87 (2005) 212108.
- [5] C.-I. Wu, C.-T. Lin, Y.-H. Chen, M.-H. Chen, Y.-J. Lu, C.-C. Wu, *Appl. Phys. Lett.* 88 (2006) 152104.
- [6] P.-C. Kao, J.-H. Lin, J.-Y. Wang, C.-H. Yang, S.-H. Chen, *Synth. Met.* 160 (2010) 1749.
- [7] P.-C. Kao, J.-Y. Wang, J.-H. Lin, S.-H. Chen, *J. Electrochem. Soc.* 157 (2010) J135.
- [8] L. Duan, Q. Liu, Y. Li, Y. Gao, G. Zhang, D. Zhang, L. Wang, Y. Qiu, *J. Phys. Chem. C* 113 (2009) 13386.
- [9] K. Xie, J. Qiao, L. Duan, Y. Li, D. Zhang, G. Dong, L. Wang, Y. Qiu, *Appl. Phys. Lett.* 93 (2008) 183302.
- [10] C. Schmitz, H.-W. Schmitz, M. Thelakkat, *Chem. Mater.* 12 (2000) 3012.
- [11] X. Sun, D.-Y. Zhou, L. Qiu, L.-S. Liao, F. Yan, *J. Phys. Chem. C* 115 (2011) 2433.
- [12] H.J. Bolink, E. Coronado, D. Repetto, M. Sessolo, *Appl. Phys. Lett.* 91 (2007) 223501.
- [13] K. Morii, M. Ishida, T. Takashima, T. Shimoda, Q. Wang, M.K. Nazeeruddin, M. Grätzel, *Appl. Phys. Lett.* 89 (2006) 183510.
- [14] N. Tokmoldin, N. Griffiths, D.D.C. Bradley, S.A. Haque, *Adv. Mater.* 21 (2009) 3475.
- [15] J. Luo, L. Xiao, Z. Chen, Q. Gong, *Appl. Phys. Lett.* 93 (2008) 133301.
- [16] Y.M. Lee, Y. Park, Y. Yi, J.W. Kim, *Appl. Phys. Lett.* 93 (2008) 123301.
- [17] Q.T. Le, L. Yan, Y. Gao, M.G. Mason, D.J. Giesen, C.W. Tang, *J. Appl. Phys.* 87 (2000) 375.
- [18] S.J. Kang, D.S. Park, S.Y. Kim, C.N. Whang, K. Jeong, S. Im, *Appl. Phys. Lett.* 81 (2002) 2581.
- [19] K. Han, Y. Yi, W.J. Song, S.W. Cho, P.E. Jeon, H. Lee, C.-N. Whang, *K. Jeong, Org. Electron.* 9 (2008) 30.
- [20] H. You, Y. Dai, Z. Zhang, D. Ma, *J. Appl. Phys.* 101 (2007) 026105.
- [21] J. Meyer, S. Hamwi, T. Bulow, H.-H. Johannes, T. Riedl, W. Kowalsky, *Appl. Phys. Lett.* 91 (2007) 113506.
- [22] L. Yan, M.G. Mason, C.W. Tang, Y. Gao, *Appl. Surf. Sci.* 175 (2001) 412.
- [23] Y.D. Jin, X.B. Ding, J. Reynaert, V.I. Arkhipov, G. Borghs, P.L. Heremans, M.V.d. Auweraer, *Org. Electron.* 5 (2004) 271.
- [24] M. Stöße, J. Staudigel, F. Steuber, J. Blässing, J. Simmerer, A. Winnacker, H. Neuner, D. Metzendorf, H.-H. Johannes, W. Kowalsky, *Synth. Met.* 111 (2000) 19.
- [25] H. Ding, Y. Gao, *Appl. Phys. Lett.* 91 (2007) 172107.
- [26] M.-Y. Chan, S.-L. Lai, K.-M. Lau, M.-K. Fung, C.-S. Lee, S.-T. Lee, *Adv. Funct. Mater.* 17 (2007) 2509.
- [27] N. Johansson, T. Osada, S. Stafstrom, W.R. Salaneck, V. Parente, D.A.d. Santos, X. Crispin, J.-L. Brédas, *J. Chem. Phys.* 111 (1999) 2157.
- [28] H. Ding, Y. Gao, *Appl. Phys. Lett.* 86 (2005) 213508.
- [29] H. Ding, Y. Gao, *Appl. Phys. Lett.* 92 (2008) 053309.
- [30] V.G. Bhide, R.H. Dani, *Physica* 27 (1961) 821.
- [31] C.H. Cheung, W.J. Song, S.K. So, *Org. Electron.* 11 (2010) 89.
- [32] S. Krause, M.B. Case, A. Scholl, E. Umbach, *New J. Phys.* 10 (2008) 085001.
- [33] A.A. Audi, P.M.A. Sherwood, *Surf. Interface Anal.* 33 (2002) 274.
- [34] P. Carreras, S. Gutmann, A. Antony, J. Bertomeu, R. Schlaf, *J. Appl. Phys.* 110 (2011) 073711.
- [35] A. Walsh, J.L.F.D. Silva, S.-H. Wei, C. Körber, A. Klein, L.F.J. Piper, A. DeMasi, K.E. Smith, G. Panaccione, P. Torelli, D.J. Payne, A. Bourlange, R.G. Egdell, *Phys. Rev. Lett.* 100 (2008) 167402.